# Diffusive Hamilton-Jacobi equations with super-quadratic growth

Alessio Porretta University of Rome *Tor Vergata* 

Singular problems associated to quasilinear equations, June 1-3, 2020

in honor of my friends Marie-Françoise and Laurent

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 in  $(0, T) \times \Omega$ 

where  $\Omega$  is a smooth bounded set in  $\mathbb{R}^N$ .

Main point of the talk:  $H(x, Du) \simeq |Du|^p$  with p > 2 $\rightsquigarrow$  beyond the natural growth.

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A short summary:

- Stationary solutions: regularity, existence & uniqueness.
  - $\rightsquigarrow$  Distributional Vs viscosity solutions

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- A short summary:
  - Stationary solutions: regularity, existence & uniqueness.
    - $\rightsquigarrow$  Distributional Vs viscosity solutions
  - Evolution of smooth solutions: when and how do we lose them ?
     → gradient blow-up, loss (and recovery !) of boundary data, ...

Remind: under natural growth conditions  $(H \le c(1 + |Du|^2))$ 

• Smooth data  $\Rightarrow$  bounded solutions are smooth

(Serrin, Trudinger, [Ladysenskaya-Uraltseva], DiBenedetto...)

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- Uniqueness of bounded weak solutions (since [Barles-Murat]...)
- Global existence for the evolution problem (either convergence or infinite time blow-up of  $||u||_{\infty}$ )

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$$\begin{cases} \lambda u - \Delta u + |Du|^{p} = f & \text{in } \Omega\\ u = 0 & \text{on } \partial \Omega \end{cases}$$

is the value function of a stochastic control problem

$$u(x) = \inf_{a \in \mathcal{A}} \mathbb{E} \left\{ \int_0^{\tau_x} \left[ c_p |a_t|^{\frac{p}{p-1}} + f(X_t) \right] e^{-\lambda t} dt \right\},$$
(1)

where  $X_t$  is a controlled process:

 $\begin{cases} dX_t = a_t \, dt + \sqrt{2} \, dB_t \, , \\ X_0 = x \; \in \Omega \, , \end{cases}$ 

 $\begin{array}{l} B_t \quad \text{a Brownian motion in } \mathbb{R}^N \\ \{a_t\}_{t\geq 0} \quad \text{a control process} \\ \tau_x = \text{ the exit time from } \Omega \end{array}$ 

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Important: the optimal drift would be given in feedback form by

$$a_t = a(X_t) = -p|Du(X_t)|^{p-2}Du(X_t)$$

 $\rightarrow$  when p > 2 singular drifts are less expensive.

## The stationary problem

$$-\Delta u + \lambda u + |Du|^p = f(x)$$
 with  $p > 2$ 

[Capuzzo Dolcetta-Leoni-P. '10]: viscosity solutions framework (fully nonlinear)  $\rightarrow$  extends to  $F(x, D^2u) + \lambda u + |Du|^p < f$ ,

(see also [Barles '10], [Barles-Koike-Ley-Topp '14] [Barles-Topp '15] for further extensions to state constraint, nonlocal diffusions etc...)

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\begin{array}{l} \mbox{[Dall'Aglio-P. '14]:} \\ \mbox{distributional solutions framework (divergence form)} \\ & \rightarrow \mbox{ extends to } -\mbox{div}(a(x, Du)) + \lambda \ u + |Du|^p \leq f, \\ \mbox{(similar with } m\mbox{-Laplacian and } p > m) \end{array}
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#### • Sub solutions are Hölder continuous

f bounded  $\Rightarrow$  USC bounded viscosity subsolutions are  $\frac{p-2}{p-1}$ -Hölder Proof by doubling variables & comparison ([Capuzzo Dolcetta-Leoni-P])

$$u(x) \leq u(y) + k\left(\frac{|x-y|}{d(x)^{1-\alpha}} + L|x-y|^{\alpha}\right) \qquad \alpha = \frac{p-2}{p-1}$$

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 $f \in L^m, m > \frac{N}{p} \Rightarrow$  distributional subsolutions are  $\alpha$ -Hölder with  $\alpha = \min(1 - \frac{N}{mp}, 1 - \frac{1}{p-1}).$ 

Proof by a Morrey-type estimate ([Dall'Aglio-P]):

$$\int_{B_r} |\nabla u|^p \, dx \leq K \, r^{N-\gamma} \,, \qquad \text{where } \gamma = \max(\tfrac{N}{m}, p')$$

## $-\Delta u + \lambda u + |Du|^p = f(x)$ with p > 2

• Interior Hölder regularity extends up to the boundary (independently of boundary data !)

Consequence: Hölder regularity is necessary for boundary data !

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• Global universal bounds for  $u^+$ :

$$||u^+||_{L^{\infty}(\Omega)} \leq M$$

where  $M = M(\Omega, \frac{1}{\lambda}, \|f\|_{L^m(\Omega)})$ , m > N/p.

Notice: the bound is independent of boundary values ! (cfr. [Lasry-Lions '89])

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#### $\rightsquigarrow$ loss of boundary data

Size and regularity conditions are needed to have compatibility of boundary data.

A sample result on the solvability of the Dirichlet problem

$$\begin{cases} \lambda u - \Delta u + |Du|^p = f & \text{in } \Omega, \\ u = \varphi & \text{on } \partial\Omega, \end{cases}$$
(2)

where p > 2 and, now, f is continuous,  $\lambda > 0$ .

#### Theorem (Capuzzo Dolcetta-Leoni-P. '10)

There exists a constant  $M_0 > 0$  such that if  $\varphi$  satisfies

$$|\varphi(\mathbf{x}) - \varphi(\mathbf{y})| \leq M |\mathbf{x} - \mathbf{y}|^{lpha} \qquad orall \mathbf{x}, \mathbf{y} \in \partial \Omega \,, \qquad lpha = rac{p-2}{p-1} \,.$$

with  $M < M_0$  and if  $\lambda$  inf  $\varphi \leq \inf f$ , then (2) has a unique viscosity solution  $u \in C^{0,(p-2)/(p-1)}(\overline{\Omega})$  such that  $u(x) = \varphi(x)$  for every  $x \in \partial \Omega$ .

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Alternatives:

- for Lipschitz solutions one can use [P-L. Lions '80]: there exists a W<sup>1,∞</sup> solution if and only if there exists a W<sup>1,∞</sup> sub solution ψ.
- For a general theory... → relaxed formulation of boundary conditions, viscosity solutions theory.

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Ex: 
$$u(x) = c_0(|x|^{\frac{p-2}{p-1}} - 1)$$
 satisfies, for a suitable choice of  $c_0$   
$$\begin{cases} -\Delta u + |\nabla u|^p = 0 & \text{in } \Omega\\ u \in W_0^{1,p}(\Omega) \cap C(\overline{\Omega}) \end{cases}$$

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 $\rightsquigarrow$  Typical first order effects

- the  $L^{\infty}$ -bound does not bring enough information...
- loss of boundary data

 $\rightarrow$  need of relaxed formulation of boundary conditions, viscosity solutions theory.

(uniqueness results for viscosity solutions: [Barles-Rouy-Souganidis '99], [Barles-Da Lio '04], [Barles '10]...)

# The evolution problem

We consider now the evolution problem

$$\begin{cases} u_t - \Delta u = |Du|^p & \text{in } (0, T) \times \Omega \\ u = 0 & \text{on } (0, T) \times \partial \Omega \\ u(0) = u_0 \end{cases}$$

- *p* > 2
- Ω smooth

• 
$$u_0 \in C^0(\overline{\Omega})$$
;  $u_0 = 0$  on  $\partial \Omega$ .

We discuss here the (more interesting) reaction case:

 $u_0 \geq 0 \quad (\Rightarrow u(t) \geq 0 \quad \forall t)$ 

Many other contributions in different directions (absorption problems, local theory in Lebesgue spaces, initial traces, etc...): [Alaa, Pierre, Barles, Da Lio, Ben Artzi, Souplet, Weissler, Bidaut Véron, Dao, Benachour, Laurençot, Dabuleanu,....]

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- But:  $||u(t)||_{C^1}$  may blow-up

$$\rightsquigarrow \quad \underline{\text{gradient blow-up:}} \quad \lim_{t \uparrow T^*} \|Du(t)\|_{\infty} = +\infty.$$

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$$\Rightarrow \quad \underline{\text{gradient blow-up:}} \quad \lim_{t\uparrow T^*} \|Du(t)\|_{\infty} = +\infty.$$

Previous contributions to the question of blow-up of classical solutions: see e.g. [Alikakos-Bates-Grant],[Conner-Grant], [Guo-Hu], [Fila-Lieberman], [Arrieta-Bernal Rodriguez-Souplet], [Li-Souplet], [Quittner-Souplet], [Souplet-Zhang], [Souplet-Vazquez],...

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(ii)  $u_t$  is bounded for all times ( $L^{\infty}$ -contraction):

$$\|u(t+h) - u(t)\|_{\infty} \le \|u(t_0+h) - u(t_0)\|_{\infty} \qquad \forall t > t_0$$

and is enough to take some  $t_0$  smaller than the blow-up time so

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(iii) At any time t, u(t) solves

$$-\Delta u = |Du|^p + f$$
 with  $f \in L^\infty$ 

Stationary gradient bounds [Lions '85] imply

$$|Du(t,x)| \leq C d(x)^{-\frac{1}{p-1}}$$
  $d(x) = \operatorname{dist}(x,\partial\Omega)$ 

#### Continuation after blow-up

• There exists a unique global *relaxed* viscosity solution

$$u \in C([0,\infty) imes \overline{\Omega}), \quad u_t - \Delta u = |Du|^p$$

[Barles-Da Lio '04] : existence and uniqueness of continuous viscosity solutions with relaxed boundary conditions:

for  $x \in \partial \Omega$ , if u(t,x) > 0 then  $u_t - \Delta u \leq |Du|^p$  in viscosity sense

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Rmk: *u* is locally Hölder continuous ([Cannarsa-Cardaliaguet '10], [Cardaliaguet-Silvestre '12], [Stokols-Vasseur '19])

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• The unique global (viscosity) solution *u* is the increasing limit of smooth solutions of truncated problems:

$$\begin{cases} \partial_t u_n - \Delta u_n = F_n(\nabla u_n), \\ u_n(0) = u_0, \quad u_{n|\partial\Omega} = 0, \end{cases} \quad \Rightarrow \quad u_n \uparrow u$$

where  $F_n$  has natural growth and  $F_n(\xi) \uparrow |\xi|^p$ 

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Classical argument: multiply by  $\varphi_1$ 

$$\int_{\Omega} u(t)\varphi_1 \, dx \geq \int_{\Omega} u_0 \varphi_1 \, dx + \int_0^t \int_{\Omega} |\nabla u|^p \varphi_1 \, dx ds - \lambda_1 \int_0^t \int_{\Omega} u \varphi_1 \, dx ds$$

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As long as u = 0 on the boundary, Poincaré inequality implies

$$\left(\int_{\Omega} u\varphi_1 \, dx\right)^p \leq \left(\int_{\Omega} u^k \, dx\right)^{p/k} \leq C\left(\int_{\Omega} |\nabla u|^k \, dx\right)^{p/k}$$
$$\leq \left(\int_{\Omega} |\nabla u|^p \varphi_1 \, dx\right) \underbrace{\left(\int_{\Omega} \varphi_1^{-k/(p-k)} \, dx\right)^{p/k-1}}_{\leq c}$$

by choosing  $1 < k < \frac{p}{2}$ .

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$$\int_{\Omega} u(t)\varphi_1 \, dx \geq \int_{\Omega} u_0 \varphi_1 \, dx + c_0 \int_0^t \left[ \left( \int_{\Omega} u \varphi_1 \, dx \right)^p - c_1^p \right] ds$$

and this would lead to a blow-up if  $\int_{\Omega} u_0 \varphi_1 dx$  is too large.

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and this would lead to a blow-up if  $\int_{\Omega} u_0 \varphi_1 dx$  is too large.  $\rightarrow u$  must lose the boundary condition (  $\Rightarrow$  gradient must blow-up)

#### • Long time behavior

[Benachour-Dabuleanu Hapca-Laurencot '07], [P.-Zuazua '12]

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What's more: there exists time  $T_0 \gtrsim K \|u_0\|_{\infty}$  such that

 $u(t) \in W_0^{1,\infty}(\Omega)$  and <u>is a classical  $C^1$ -solution</u> for  $t \ge T_0$  and  $\|Du(t)\|_{\infty} \le C \frac{e^{-\lambda_1 t}}{\sqrt{t}} \quad \forall t \ge T_0$ 

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So not only there is life after blow-up, but there is also a happy ending  $\rightsquigarrow$  whatever  $u_0$  is, solutions eventually become classical again and behave like the heat equation

So far, three main features appear in the qualitative behavior of the evolution problem:

(i) gradient blow-up

(blow-up rate, blow-up profile,...?)

- (ii) loss of boundary condition(when and how does it occur ?)
- (iii) recovery of boundary condition and regularization(a totally new issue ! How this smoothing property can be described?...)

Jointly with with P. Souplet, we investigated the relations between these 3 issues.

## 1. Blow-up rate (and regularization rate)

In [P-Souplet '19] we give some estimates for the rate of blow-up and regularization:

$$T^* :=$$
 blow-up time  $T^r :=$  regularization time

The regularization time  $T^r = T^r(u_0)$  means the (ultimate) time after which the solution remains classical for ever:

$$T^r(u_0):=\inf\bigl\{\tau>T^*(u_0);\ u(t,\cdot)\in C^1_0(\overline\Omega) \text{ for all } t>\tau\bigr\}\in [T^*,\infty).$$

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In both cases, we show that the minimal rate has order 1/(p-2):

 $\text{if blow-up time } \mathcal{T}^* < \infty \text{, then} \quad \| \nabla u(t) \|_\infty \geq \frac{\mathcal{C}}{(\mathcal{T}^* - t)^{1/(p-2)}}, \quad t \uparrow \mathcal{T}^*$ 

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and similarly at the regularization time:

$$\|
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Remarks:

 The blow-up rate (T\* - t)<sup>-1/(p-2)</sup> is known to be optimal for time-increasing solutions (1d or radial, see [Conner-Grant], [Guo-Hu],[Attouchi-Souplet]).

However, faster blow-up rates can occur !! (see later....)

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However, faster blow-up rates can occur !! (see later....)

• This is not the self-similar rate of the equation. The self-similar scaling would lead to a rate of order 1/(p-1).

$$u_{\lambda}(x, y, t) := \lambda^{m} u(\lambda x, \lambda y, \lambda^{2} t) \quad \text{with } m = (2 - p)/(p - 1),$$
  
$$\rightsquigarrow \quad \nabla u_{\lambda} = \lambda^{1/(p-1)} \nabla u(\lambda x, \lambda y, \lambda^{2} t).$$

The self-similar rate is actually the rate of the spatial blow-up of normal derivative (see also [Filippucci-Pucci-Souplet])

~→ for single point blow-up, in 2 – d we proved ([P-Souplet '17]) that the blow-up profile is strongly anisotropic: 1/(p-2) for the time rate 1/(p-1) of the normal spatial rate 2/(p-2) for the tangential spatial rate

## 2. Loss of boundary conditions

In [P.-Souplet '17], we show that loss of boundary condition may or may not occur, strongly depending on the initial data.

• there exist initial data for which the loss of boundary conditions occurs everywhere on  $\partial \Omega$ 

(see also [Quaas-Rodriguez] in the setting of viscosity solutions)

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• one can find solutions for which the loss of boundary conditions occurs essentially only on a prescribed open subset of  $\partial\Omega$ 

 $\forall \omega$  open set of  $\mathbb{R}^n$ , and  $\forall \varepsilon > 0$ , there exists  $u_0$ :

$$\omega \cap \partial \Omega \subset \Sigma_u \subset [\omega + B_{\varepsilon}(0)] \cap \partial \Omega$$

where  $\Sigma_u$  is the set where the boundary condition is lost.

In other words, one can prepare the initial datum  $u_0$  so that loss of boundary condition occurs at his/her favorite subset of  $\partial\Omega$ 

#### • gradient blowup may occur without loss of boundary conditions

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[P-Souplet] in 1-d, and [Filippucci-Pucci-Souplet] in more generality

Actually, for fixed  $u_0 \ge 0$ , this happens at the critical value

$$ar{\lambda} := \inf\{\lambda > 0 \ : \ T^*(\lambda u_0) < \infty\}$$

If  $\lambda u_0 \rightsquigarrow u_\lambda$ , then

- $u_{\lambda}$  blows-up if and only if  $\lambda \geq \overline{\lambda}$
- $u_{\lambda}$  loses the boundary value if and only if  $\lambda > \bar{\lambda}$

 $\rightsquigarrow u_{\bar{\lambda}}$  is a *threshold* between global in time smooth solutions and solutions with loss of boundary condition

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This distinguishes between *minimal blow-up solutions* and *non minimal* ones:

*u* is a minimal blow-up solution if, for any  $v_0 \leq u_0$ ,

the corresponding solution v is smooth for all times

Be careful: minimal blow-up solutions may have faster blow-up rates!

# **3.** The one-dimensional case: a precise picture of loss and recovery of boundary conditions.

Take initial data  $u_0$  which satisfy

$$u_0$$
 is symmetric w.r.t.  $x = \frac{1}{2}$ ,  $u'_0 \ge 0$  on  $[0, \frac{1}{2}]$ ,  
 $u_0(0) = u''_0(0) + {u'_0}^p(0) = 0$ 

there exists 
$$a \in (0,1)$$
 such that  $u_0'' + u_0'{}^p \begin{cases} > 0 & \text{on } [0,a) \\ < 0 & \text{on } (a, \frac{1}{2}]. \end{cases}$  (3)

Set as before

 $T^* =$ blow-up time  $T^r =$ regularization time

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[P-Souplet '19]: a detailed picture of minimal and non minimal solutions in 1-d

1. *u* is a minimal blow-up solution if and only if it does not lose the boundary condition. Moreover:

(i) There is instantaneous and permanent regularization at  $T^*$ , *u* becomes immediately smooth once and for ever after  $T^*$ .

(ii) both the blow-up and the regularization rate are faster than the minimal rate:

$$\begin{split} (T^*-t)^{1/(p-2)}\|u_x(t)\|_{\infty} &\to \infty, \quad t\uparrow T^* \\ (t-T^*)^{1/(p-2)}\|u_x(t)\|_{\infty} &\to \infty, \quad t\downarrow T^* \end{split}$$

Rmk: for a class of non monotone in time solutions, [Attouchi-Souplet] recently proved a blow-up rate of 2/(p-2) !...

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#### (i) The blow-up and regularization rates are both minimal:

$$c_1(\mathcal{T}^*-t)^{-1/(p-2)} \leq \|u_x(t)\|_\infty \leq c_2(\mathcal{T}^*-t)^{-1/(p-2)}, \quad ext{ as } t o \mathcal{T}^*_-,$$

$$c_3(t-T^r)^{-1/(p-2)} \leq \|u_x(t)\|_\infty \leq c_4(t-T^r)^{-1/(p-2)}, \quad ext{ as } t o T^r_+,$$

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(ii) u linearly detaches from the boundary condition and linearly reconnects:

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(iii) There is immediate and permanent regularization after the recovery of boundary conditions

$$u(t, 0) > 0$$
 for all  $t \in (T^*, T^r)$ 

then u becomes a classical solution again and for ever

We also describe the life of the solution in the time interval when the boundary condition is relaxed. Actually, we have

(iii) In the interval  $[T^*, T^r]$ , the solution behaves like a shifted copy of the singular stationary profile:

$$u(t,x) \simeq u(t,0) + c_p x^{-1/(p-1)}$$
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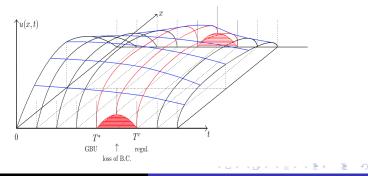
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Moreover

u(t,0) is  $C^1$  in  $[T^*, T^r]$ , admits a unique max. at some  $T_m \in (T^*, T^r)$ , it is increasing on  $[T^*, T_m]$  and decreasing on  $[T_m, T^r]$ 



•  $u_t$  satisfies a linear equation  $z_t - z_{xx} = bz_x$ ,  $b = p|u_x|^{p-2}u_x$ , so up to  $t = T^*$  one can apply the zero number property ([Matano], [Angenent])

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 $\rightsquigarrow$  there are (roughly speaking) only two cases:

(i) either the zero of  $u_t$  collapses at the boundary as  $t \uparrow T^*$ Then  $u_t \leq 0$  at  $t = T^*$  (and later), and the boundary condition can not be lost (minimal solution)

(ii) or  $u_t$  remains positive near x = 0 at  $t = T^*$ , then  $u_t(t, 0)$  jumps at  $t = T^*$  and u leaves the boundary condition and becomes positive on  $\partial \Omega$  (non minimal solution)

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Important: Recent results by [Mizoguchi-Souplet] show that the behavior can be much more weird for other initial data: gradient blow-up and loss/recovery of boundary condition may happen several times !!

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Conclusions, comments,...

- Perturbing the Laplace (or the heat) equation with super quadratic coercive first order terms lead to a surprisingly rich hybrid model of a second order equation which shares properties of first order equations.
- The stationary equation already shows the possible competition between oscillations (due to the Laplacian) and coercivity (given by the Hamiltonian).

It gives much more rigidity compared to a second order equation: the coercive term may prevent too singular oscillations and gives prescriptions on size and regularity of solutions up the boundary Conclusions, comments,...

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Once more, the competition between the possible formation of singularity (gradient blow-up) and the stabilizing effects (due to the maximum principle) leaves room for a rich variety of phenomena: anisotropic blow-up rates, profiles, etc...

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# Thanks for the attention !

# Happy Birthday to Marie-Françoise and Laurent !

Wishing you and us all that for your next celebration we will have recovered our boundary conditions...:

- coffee breaks
- dinners to share
- random walks in stranger countries
- personal discussions
- claps after the talks...:)
- ....